

# Kinetic Study of Hydrogen Evolution Reaction over Strained MoS<sub>2</sub> with Sulfur Vacancies Using Scanning Electrochemical Microscopy

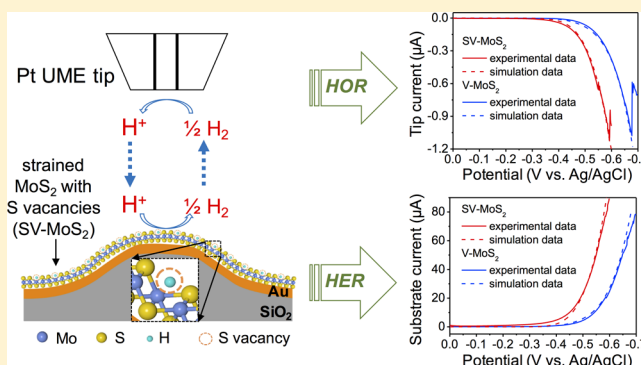
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## Supporting Information

**ABSTRACT:** Molybdenum disulfide (MoS<sub>2</sub>), with its active edge sites, is a proposed alternative to platinum for catalyzing the hydrogen evolution reaction (HER). Recently, the inert basal plane of MoS<sub>2</sub> was successfully activated and optimized with excellent intrinsic HER activity by creating and further straining sulfur (S) vacancies. Nevertheless, little is known about the HER kinetics of those S vacancies and the additional effects from elastic tensile strain. Herein, scanning electrochemical microscopy was used to determine the HER kinetic data for both unstrained S vacancies (formal potential  $E_v^0 = -0.53$  V<sub>Ag/AgCl</sub>, electron-transfer coefficient  $\alpha_v = 0.4$ , electron-transfer rate constant  $k_v^0 = 2.3 \times 10^{-4}$  cm/s) and strained S vacancies ( $E_{sv}^0 = -0.53$  V<sub>Ag/AgCl</sub>,  $\alpha_{sv} = 0.4$ ,  $k_{sv}^0 = 1.0 \times 10^{-3}$  cm/s) on the basal plane of MoS<sub>2</sub> monolayers, and the strained S vacancy has an electron-transfer rate 4 times higher than that of the unstrained S vacancy. This study provides a general platform for measuring the kinetics of two-dimensional material-based catalysts.



## INTRODUCTION

Molybdenum disulfide (MoS<sub>2</sub>) has been extensively studied as an attractive catalyst for the hydrogen evolution reaction (HER) for about three decades.<sup>1–4</sup> The edge sites of MoS<sub>2</sub> were first theoretically proposed as active sites in 2005<sup>5</sup> and then were experimentally confirmed in 2007.<sup>6</sup> Since then, efforts to optimize MoS<sub>2</sub> for HER have focused on maximizing the density of edge sites.<sup>7–9</sup> Very recently, the basal plane of MoS<sub>2</sub>, which is long thought to be inert for HER, was successfully activated by creating sulfur (S) vacancies that generated new gap states near the Fermi level for hydrogen binding<sup>10,11</sup> and further enhanced by elastically straining the S vacancies that moved the new gap states closer to the Fermi level.<sup>11</sup> So far, all the understanding of the HER activity of those unstrained or strained S vacancies comes from thermodynamic consideration of their hydrogen adsorption energy,<sup>11</sup> and little is known about the kinetics of those S vacancies on the basal plane of MoS<sub>2</sub>.

In this report, we employ scanning electrochemical microscopy (SECM)<sup>12</sup> equipped with a platinum (Pt) ultramicroelectrode (UME) to determine the kinetic information on both unstrained and strained S vacancies on the basal plane of MoS<sub>2</sub> monolayers. SECM has been previously applied to study the HER kinetics of metallic catalysts, such as Pt UME,<sup>13</sup> manganese (Mn) electrodes,<sup>14</sup> Pt and gold (Au) thin layer cells,<sup>15</sup> and metal nanoparticles.<sup>16–18</sup> These studies

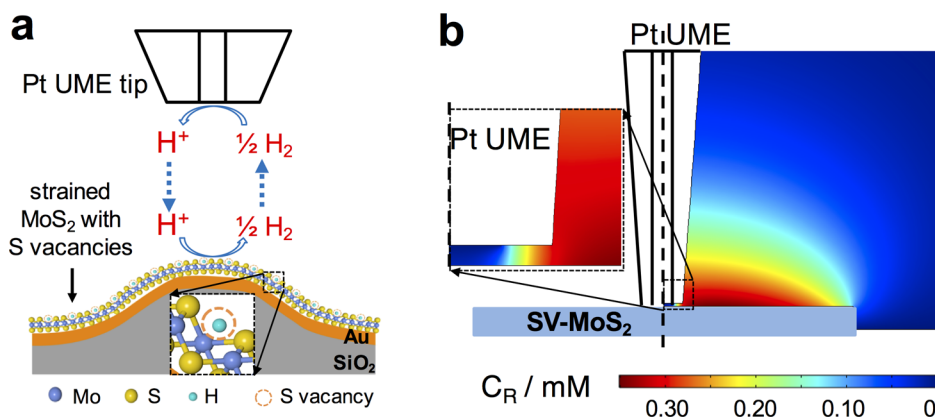
successfully determined the apparent rate constants and electron-transfer coefficients of the HER on Pt, palladium (Pd), and Mn electrodes.<sup>13,14,16,17</sup> More importantly, these studies established the methodology to quantify the HER kinetics, including simplifying the HER to be a quasi-reversible one-step reaction, using a thin layer cell formalism of the HER in the SECM theoretical equations<sup>19</sup> and describing the rates of the elementary steps and coverage of adsorbed hydrogen atoms on metallic electrodes based on the Volmer–Heyrovsky–Tafel mechanism.<sup>15</sup> Herein, we leverage the existing framework of the HER study (using SECM to investigate heterogeneous electrochemical reactions in the energy materials)<sup>20–22</sup> and adapt it to study the HER kinetics of the two-dimensional (2-D) catalyst of MoS<sub>2</sub> monolayers.

## EXPERIMENTAL SECTION

**Sample Fabrication.** Continuous 2H-MoS<sub>2</sub> monolayers were grown in a homemade tube furnace.<sup>11</sup> The as-grown MoS<sub>2</sub> monolayers were transferred onto a silicon dioxide (SiO<sub>2</sub>) substrate capped by a gold (90 nm)/titanium (10 nm) film using the poly(methyl methacrylate) (PMMA)-assisted transfer process. Briefly, a PMMA film was coated onto the as-grown MoS<sub>2</sub> on a SiO<sub>2</sub> substrate. Then the sample was soaked in potassium hydroxide solution to etch the surface

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**Figure 1.** Methodology for quantifying the HER kinetics of strained/unstrained S vacancies on the basal plane of MoS<sub>2</sub> monolayers. (a) Schematic of the SECM measurement setup that operates in the SG–TC mode. The top trapezoid represents a Pt UME tip, and the bottom strained MoS<sub>2</sub> monolayer with S vacancies functions as the working electrode. H<sub>2</sub> is generated at the MoS<sub>2</sub> substrate through HER and subsequently collected and oxidized at the Pt UME tip through HOR. (b) Schematic of the complementary COMSOL simulation for the spatial distribution of hydrogen concentration ( $C_R$ /mM). The left side shows the tip and substrate 2-D configuration used in the simulation. The tip is a 25  $\mu\text{m}$  diameter Pt wire in a glass shell (the ratio of glass to metal radius  $R_G = 2$ ). The SV-MoS<sub>2</sub> substrate has a diameter of about 0.5 mm. The right side shows a representative calculated  $C_R$  distribution (condition: SV-MoS<sub>2</sub> potential at  $-0.6$  V and Pt UME tip potential at 0 V in 0.1 M HClO<sub>4</sub>). Left inset: zoomed-in view of  $C_R$  near the tip–substrate gap. Bottom inset: color bar of  $C_R$ .

SiO<sub>2</sub> layer and to lift off the PMMA/MoS<sub>2</sub> bilayer film (see Supporting Note 1 for more details). Afterward, the PMMA/MoS<sub>2</sub> bilayer film was cleaned and floated on a deionized water surface and then scooped by the target substrate. In the center of the SiO<sub>2</sub> substrate, an array of SiO<sub>2</sub> nanopillars (Supporting Information) was obtained via dry etching of the SiO<sub>2</sub> film with an electron-beam resist as the etching mask, which was patterned by electron-beam lithography. Columns were defined with a negative-tone electron-beam resist (MaN-2403). The pitch size was 650 nm, and the diameter and height of the nanopillars were 350 and 80 nm, respectively. The gold nanopillars were formed after metallization of SiO<sub>2</sub> nanopillars. After thorough PMMA removal by sequential soaking of samples in acetone (60 °C for 10 min) and chloroform (60 °C for 1 h), the sample was strained by capillary force (see Supporting Note 1).<sup>23</sup> To create S vacancies, the cleaned MoS<sub>2</sub> samples were exposed to a cold and mild argon plasma (Supporting Note 1).<sup>24</sup> The MoS<sub>2</sub> film on the flat area of the substrate (outside the nanopillar array area) was used as MoS<sub>2</sub> with S vacancies (V-MoS<sub>2</sub>), and that inside the nanopillar array area was used as strained MoS<sub>2</sub> with S vacancies (SV-MoS<sub>2</sub>). The same PMMA-assisted transfer method was used to prepare the transmission electron microscopy (TEM) sample for S vacancy characterization by transferring the as-grown MoS<sub>2</sub> monolayers onto a TEM grid (quantifoil TEM grid, Ted Pella, Inc.), followed by PMMA removal through sequential acetone and chloroform soaking processes. Then the TEM sample was treated with mild Ar plasma to create S vacancies (Supporting Note 1).

**Sample Characterization.** Elastic strain in the sample was quantified from the well-established strain-Raman shift relation. Raman modes ( $E_{2g}^1$  and  $A_{1g}$ ) of monolayer MoS<sub>2</sub> show red shifts upon uniaxial tensile strain with ratios of 2.2 and 0.5  $\text{cm}^{-1}/\%$  for  $E_{2g}^1$  and  $A_{1g}$  modes, respectively.<sup>23,25</sup> The magnitude of the elastic strain was thus calculated by the amount of red shifts of Raman modes (Figure S1). Aberration-corrected, monochromated TEM (ACTEM) imaging was performed using a FEI spherical aberration (Cs) image-corrected 80–300 Titan environmental (scanning) TEM operated at 80 kV under negative Cs imaging (NCSI) conditions,<sup>26,27</sup> wherein the spherical aberration coefficient of the objective lens was adjusted to approximately  $-15$   $\mu\text{m}$ , and images were acquired using a slight overfocus setting.

**MoS<sub>2</sub> Working Electrode Preparation.** The MoS<sub>2</sub> working electrodes were defined by masking the MoS<sub>2</sub> substrate using plastic tape with two hole openings. One hole was located on the unstrained MoS<sub>2</sub> area (V-MoS<sub>2</sub>), and the other was on the nearby strained MoS<sub>2</sub>

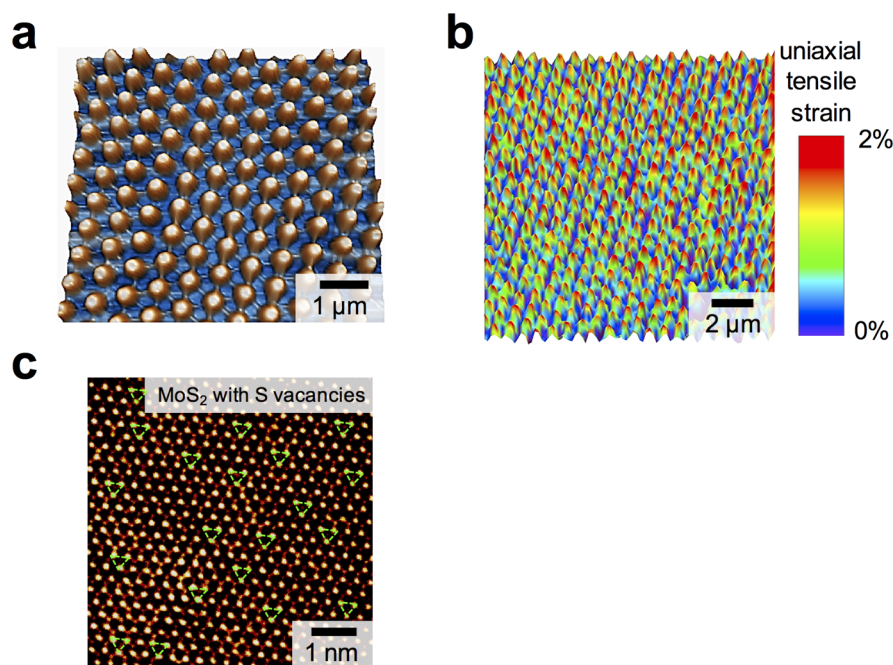
area (SV-MoS<sub>2</sub>). Only the areas within the two holes were exposed to the electrolyte (Supporting Note 2 and Figure S2a).

**Electrochemical Measurements.** A 920C SECM (CH Instruments, Austin, TX) setup (Figure S2b) was employed for the electrochemical measurements. The MoS<sub>2</sub> samples on a Au/Ti-coated SiO<sub>2</sub>/Si wafer were placed between an acrylic base and a PTFE cell with a 3.8 mm diameter orifice in the center. The samples were then connected to the potentiostat through a copper tape that was stuck to the bare Au part of the SiO<sub>2</sub>/Si wafer. An O-ring with a 1.8 mm inner diameter was used to ensure a leakage-free connection. A 25  $\mu\text{m}$  diameter Pt UME served as the tip, which was mounted facing downward to the substrate electrode. A Pt wire was used as a counter electrode, and a Ag/AgCl (1 M) electrode (CH Instruments, Austin, TX) acted as a reference electrode. All potentials shown in this work were relative to the Ag/AgCl (1 M) reference electrode unless stated otherwise. For the SECM experiments, a sealed glovebag was used to keep the system free of oxygen by purging with argon (Ar). All SECM measurements were performed in 0.1 M HClO<sub>4</sub> aqueous solution after deaeration via Ar bubbling.

The V-MoS<sub>2</sub> and SV-MoS<sub>2</sub> substrate electrodes were studied successively by being individually placed within the O-ring area. Two kinds of substrate generation–tip collection (SG–TC) modes were applied. First, the MoS<sub>2</sub> substrate potential was swept from 0 V to  $-0.7$  V at a scan rate of 5 mV/s to generate hydrogen continuously, and the hydrogen was then collected at the Pt UME tip. Second, the MoS<sub>2</sub> substrate potential was switched between the open circuit potential (OCP) and the negative potentials for HER by dual-potential pulses. The pulse width was 0.5 s with five replicates. Also, the tip potential was held at 0 V as an amperometric sensor for collecting and oxidizing the hydrogen.

The measurements in Figure S6 were performed by a three-electrode electrochemical compression cell connected to a Gamry 600 potentiostat. The MoS<sub>2</sub> substrate served as the working electrode; a Pt wire acted as the counter electrode, and a Ag/AgCl electrode was used as the reference electrode. Sulfuric acid (0.3 M) with a pH value of 0.2 was used as the electrolyte.

**Multiphysics Simulation.** For SECM SG–TC experiments, numerical simulations were performed using COMSOL Multiphysics software v4.2 (COMSOL, Inc., Burlington, MA). A 2-D axial symmetry model was employed, where a 12.5  $\mu\text{m}$  radius Pt UME ( $a = 12.5$   $\mu\text{m}$ ,  $R_G = 2$ ) and a 250  $\mu\text{m}$  radius MoS<sub>2</sub> substrate ( $b = 250$   $\mu\text{m}$ ) were vertically aligned in a cylinder with a radius of 500  $\mu\text{m}$  and a height of 500  $\mu\text{m}$ . The distance between the tip and substrate was 5  $\mu\text{m}$  for V-MoS<sub>2</sub> and 4  $\mu\text{m}$  for SV-MoS<sub>2</sub> modeling. Since the scan rate



**Figure 2.** Quantification of the strain magnitude and S vacancy concentration in MoS<sub>2</sub> monolayers. (a) AFM image of a typical strained monolayer MoS<sub>2</sub> sample on patterned Au nanopillars. Scale bar is 1 μm. (b) Color-coded strain distribution in a strain-textured monolayer MoS<sub>2</sub> (14 × 14 μm<sup>2</sup>) obtained from scanning micro-Raman spectroscopy. The peak strain magnitude on the top of nanopillars is about 2% (uniaxial tensile strain). (c) AC-TEM image of a ~6.5 × 6.5 nm<sup>2</sup> MoS<sub>2</sub> monolayer with about ~13% S vacancies. The biggest and brightest dots are Mo atoms, and the small bright dots are S atoms. Small dim dots are mono-S vacancies, some of which have been labeled by dashed triangles. Scale bar is 1 nm.

of the SG–TC experiment was sufficiently slow (5 mV/s), we used the steady-state solver to perform the simulation for the potential sweep experiments. The kinetic parameters of the HER for V-MoS<sub>2</sub> and SV-MoS<sub>2</sub>, including the formal potential  $E^0$ , electron-transfer coefficient  $\alpha$ , and the rate constants  $k^0$ , were determined by finding the best fit between experimental and simulated linear sweep voltammograms for both the tip and the substrate (see Supporting Note 3 for more details).

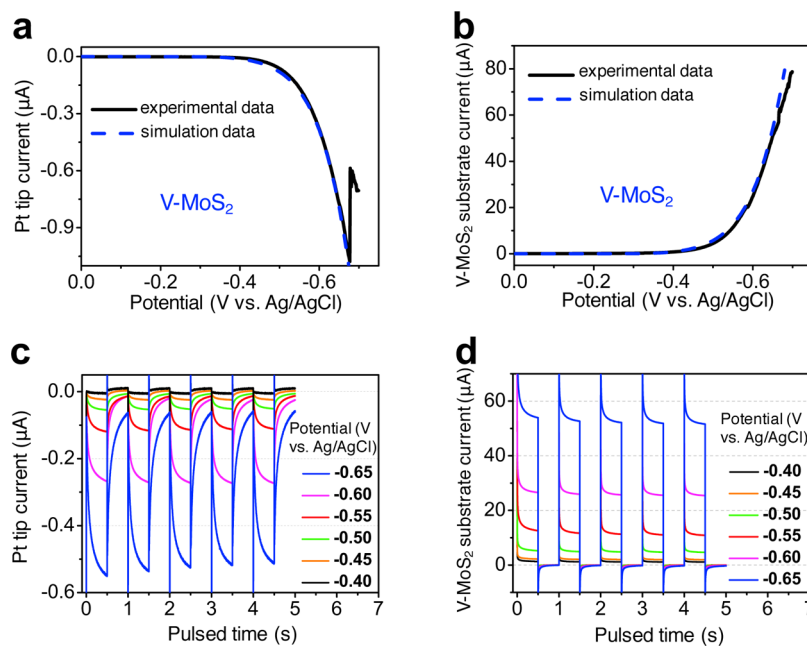
## RESULTS AND DISCUSSION

The SECM measurement setup is schematically illustrated in Figure 1a, and it operates in SG–TC mode.<sup>28,29</sup> The monolayer MoS<sub>2</sub> supported on Au is the substrate electrode (i.e., working electrode) that is negatively biased to evolve H<sub>2</sub> (i.e., HER:  $\text{H}^+ + \text{e}^- \rightarrow \frac{1}{2}\text{H}_2$ ), while a Pt UME tip is held at constant 0 V vs Ag/AgCl (1 M) reference electrode and measures the current generated from the hydrogen oxidation reaction (i.e., HOR:  $\frac{1}{2}\text{H}_2 \rightarrow \text{H}^+ + \text{e}^-$ ). Since the HOR current collected by the Pt UME tip is controlled by the diffusion of H<sub>2</sub> generated on the MoS<sub>2</sub> monolayer electrode, it excludes potential current contributions from other side reactions (e.g., reduction of growth precursor, molybdenum trioxide). As such, the SG–TC mode allows us to accurately determine the electron-transfer kinetics of the HER of the MoS<sub>2</sub> monolayer electrode. Complementary finite difference simulation is carried out to calculate the time-dependent concentration profiles of hydrogen and the proton in the SG–TC mode of SECM (Figure 1b). Since there is no recognized model for the HER kinetics of MoS<sub>2</sub>, we assume that it follows a general one-electron-transfer reaction model with a Butler–Volmer formalism. All kinetic parameters were determined by identifying the best fit between the theoretically simulated and experimentally measured linear sweep voltammogram (LSV) curves. It is worth noting that the work described here concerns the relative rates of the HER

reaction, thus the model selected to fit the data is of secondary importance.

For the SECM measurement, the working electrode consists of either unstrained (on flat substrates, V-MoS<sub>2</sub>) or strained (on nanopillar arrays, Figure 1a shows a single nanopillar for a clear illustration, SV-MoS<sub>2</sub>) MoS<sub>2</sub> monolayers with S vacancies. The atomic force microscopy (AFM) image in Figure 2a shows that the MoS<sub>2</sub> monolayers conformally and uniformly coat the nanopillar array. As such, MoS<sub>2</sub> monolayers are strained with larger strain on top of the nanopillars and smaller strain between nanopillars.<sup>23</sup> The magnitude and spatial distribution of the strain (Figure 2b) were determined by Raman spectroscopy,<sup>25,30</sup> and the peak strain magnitude is about 1% biaxial tensile strain, or equivalently 2% uniaxial tensile strain (see Figure S1). The S vacancies on the basal plane of MoS<sub>2</sub> monolayers are clearly shown in the AC-TEM image (Figure 2c). In this image, the brightest spots correspond to Mo atoms and the lower intensity spots are the S atoms. The highlighted areas (dashed triangles) in Figure 2c correspond to the location of S atoms, but they exhibit dimmer or even absent brightness, indicating that one or two (top and bottom atomic planes) S atoms in the MoS<sub>2</sub> monolayer are missing. By counting the number of those S atoms with dimmer or absent brightness in the AC-TEM image, we estimate that about 13% S atoms have been removed by the Ar plasma treatment.

Next, we ensure that the Pt UME tip is placed in the hydrogen diffusion layer above the MoS<sub>2</sub> working electrodes for the SECM measurements, so that the HOR current collected by the Pt UME tip is controlled by the diffusion of H<sub>2</sub> generated on the MoS<sub>2</sub> working electrodes. Specifically, the Pt UME tip was first electrochemically activated in Ar-saturated 0.1 M HClO<sub>4</sub> by cycling the applied potential between –0.2 and +1.0 V at a scan rate of 0.1 V/s for 20 cycles. Afterward, the Pt UME tip was held at a fixed potential at –0.35 V, where



**Figure 3.** Experimental SECM measurement results of V-MoS<sub>2</sub> and the corresponding simulation results. (a) Pt tip and (b) V-MoS<sub>2</sub> substrate voltammograms in 0.1 M HClO<sub>4</sub> electrolyte and the corresponding simulation results (blue dashed lines). The separation between the tip and substrate electrode was  $\sim 5 \mu\text{m}$ . Tip potential was held at 0 V, and the substrate potential was swept from 0 V to  $-0.7$  V at a scan rate of 5 mV/s. Transient and oscillating (c) Pt tip and (d) V-MoS<sub>2</sub> substrate current according to the double-pulse potential applied to the substrate. Tip potential was held at 0 V, and the substrate potential was switched between open circuit potential and the potentials listed in the legend to enable MoS<sub>2</sub> for HER. The pulse width was 0.5 s with five replicates.

proton reduction occurred, and the tip approached the MoS<sub>2</sub> substrate in the normal  $z$  direction using the “negative feedback mode”,<sup>31,32</sup> while the MoS<sub>2</sub> substrate electrode was disconnected from the SECM (see Figure S2 for the setup). The tip approach curves for both V-MoS<sub>2</sub> and SV-MoS<sub>2</sub> samples were experimentally recorded and calculated with the negative feedback mode (Figure S3a,b). Excellent agreement between the measured and calculated tip approach curves confirms that the Pt tip approached the working substrate in the negative feedback mode as designed. As the Pt tip became closer to the MoS<sub>2</sub> substrate, its current decreased because the diffusion of the proton to the tip was hindered by the physical obstruction of the substrate.<sup>33</sup> When the Pt tip current drops to about 60% of its initial value ( $I/I_\infty \approx 0.6$ ), it normally indicates that the Pt tip is in the hydrogen diffusion layer and further approach to the substrate may crash the Pt tip. As such, the position for  $I/I_\infty \approx 0.6$  was used to obtain the final tip–substrate distance, which was determined to be about 5  $\mu\text{m}$  for V-MoS<sub>2</sub> and 4  $\mu\text{m}$  for SV-MoS<sub>2</sub>. Once the tip was positioned at the set distances over the MoS<sub>2</sub> substrates, the MoS<sub>2</sub> substrates were electrochemically activated and stabilized by cycling between 0 V and  $-0.5$  V in 0.1 M HClO<sub>4</sub> at a scan rate of 20 mV/s for about five cycles. After stabilization, the LSV measurements were performed by sweeping the applied potential on MoS<sub>2</sub> from 0 V to  $-0.7$  V at a scan rate of 5 mV/s to generate H<sub>2</sub> while holding the Pt UME tip potential at 0 V to oxidize the generated H<sub>2</sub> into protons.

Figure 3 shows representative SECM LSV plots of the Pt tip current (Figure 3a) and substrate current (Figure 3b) for which the substrate is the unstrained MoS<sub>2</sub> monolayer with S vacancies (V-MoS<sub>2</sub>). The abrupt decrease in tip current at  $-0.68$  V in Figure 3a was caused by the formation of hydrogen bubbles, which crossed the tip–substrate gap and made contact with the Pt UME. The overall current collection efficiency, that

is, the ratio between the tip current and the substrate current, is small at about 1.37%, which is reasonable because the tip size is much smaller than the size of the substrate. To obtain the HER kinetic parameters for the V-MoS<sub>2</sub> substrate, 2-D axial symmetric COMSOL simulations were carried out to calculate the time-dependent H<sup>+</sup> (with concentration  $C_O$ ) and H<sub>2</sub> (with concentration  $C_R$ ) distribution (Figure 1b; see Supporting Information for details).<sup>34,35</sup> For the simulation, the forward ( $K_{f_v}$ ) and backward ( $K_{b_v}$ ) HER reaction rates on the V-MoS<sub>2</sub> working electrode were calculated using the following equations:

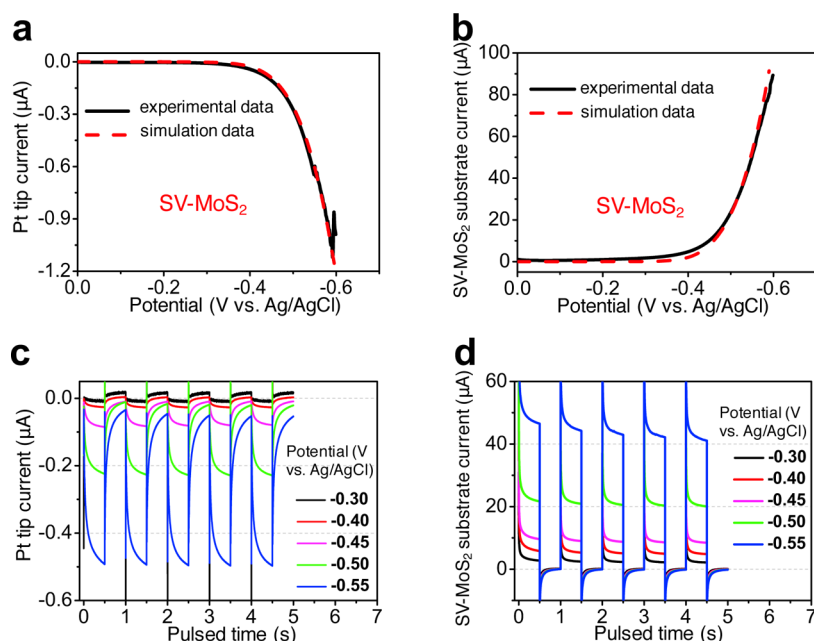
$$K_{f_v} = k_v^0 \times e^{[-\alpha_v n f (E_{\text{sub}} - E_v^0)]}$$

$$K_{b_v} = k_v^0 \times e^{[(1-\alpha_v) n f (E_{\text{sub}} - E_v^0)]}$$

where  $f = F/RT$  [ $\sim 38.92/\text{V}$ , given Faraday’s constant  $F = 9.6 \times 10^4$  C/mol, the ideal gas constant  $R = 8.3$  J/(K·mol), and temperature  $T = 298$  K] and  $E_v^0$ ,  $\alpha_v$ , and  $k_v^0$  are the corresponding formal potential, electron-transfer coefficient, and electron-transfer rate constant to be determined. For the HOR on the Pt tip, its forward and backward reaction rates were calculated using the reported electron-transfer kinetic parameters of  $E_t^0 = -0.40$  V,  $\alpha_t = 0.5$ , and  $k_t^0 = 0.42$  cm/s.<sup>13,34</sup> Since the SECM operated in the diffusion-controlled regime, the calculated  $C_O$  and  $C_R$  profiles were then used to calculate the substrate HER and Pt tip HOR current by using Fick’s first law of diffusion:

$$\frac{i_{\text{HER}}}{nFA} = D_O \frac{\partial C_O}{\partial z}, \quad \frac{i_{\text{HOR}}}{nFA} = D_R \frac{\partial C_R}{\partial z}$$

where  $i$  is the current,  $n$  is the electron number,  $A$  is the electrode area, and  $D_O$  and  $D_R$  are the proton and H<sub>2</sub> diffusion coefficients, respectively.<sup>34,35</sup> By identifying the best fit between



**Figure 4.** Experimental SECM measurement results of SV-MoS<sub>2</sub> and the corresponding simulation results. (a) Pt tip and (b) SV-MoS<sub>2</sub> substrate voltammograms in 0.1 M HClO<sub>4</sub> electrolyte and the corresponding simulation results (red dashed lines). The separation between the tip and substrate electrode was  $\sim 4 \mu\text{m}$ . Tip potential was held at 0 V, and the substrate potential was swept from 0 V to  $-0.6$  V at a scan rate of 5 mV/s. Transient and oscillating (c) Pt tip and (d) SV-MoS<sub>2</sub> substrate current according to the double-pulse potential applied to the substrate. Tip potential was held at 0 V, and the substrate potential was switched between open circuit potential and the potentials listed in the legend to enable MoS<sub>2</sub> for HER. The pulse width was 0.5 s with five replicates.

the calculated currents (Figure 3a,b, dashed lines) and the measured currents (Figure 3a,b, solid lines), the HER kinetic data ( $E_v^0 = -0.53$  V,  $\alpha_v = 0.4$ ,  $k_v^0 = 2.3 \times 10^{-4}$  cm/s) were determined for V-MoS<sub>2</sub>.

Since the HER would not reach the steady state in 0.1 M HClO<sub>4</sub> acid during our SECM measurement, we further conducted transient response measurements to verify the above kinetic parameters extracted from the LSV curves. For the transient response study,<sup>36</sup> we applied a pulse potential to the MoS<sub>2</sub> substrate electrode between the OCP and various negative potentials with a pulse width of 0.5 s for five replicates. The applied potential to the Pt tip was held at 0 V, which was much more positive than the HOR formal potential ( $-0.40$  V), so that the HOR proceeded at a diffusion-controlled rate on the Pt tip, which ensured the validity of extracting the HER kinetic parameters from the tip current, as discussed earlier. The corresponding transient Pt tip and MoS<sub>2</sub> substrate currents are plotted in Figure 3c,d, respectively, and both show similar oscillating behavior. First, the Pt tip current (Figure 3c) is nearly 0 A when a MoS<sub>2</sub> substrate is at the OCP (without any HER activity), which shows no hysteresis for H<sub>2</sub> collection on the Pt tip. Second, when more negative potential is applied to MoS<sub>2</sub>, it generates a larger HER substrate current and larger Pt tip HOR current, which verifies that the Pt tip HOR current is indeed controlled by the H<sub>2</sub> diffusion flux. Moreover, it is noted that both the tip and the V-MoS<sub>2</sub> substrate currents agree with the electron-transfer kinetics obtained from Figure 3a,b, which also suggests that the one-electron reaction model with a Butler–Volmer equation provides a reasonable formalism to treat the complex HER kinetics on V-MoS<sub>2</sub>.

Next, we applied similar SECM measurement steps (as discussed above for V-MoS<sub>2</sub>) to the SV-MoS<sub>2</sub> sample with a current collection efficiency of  $\sim 1.26\%$ . The corresponding kinetic data for SV-MoS<sub>2</sub> were determined to be  $E_{sv}^0 = -0.53$  V,

$\alpha_{sv} = 0.4$ , and  $k_{sv}^0 = 1.0 \times 10^{-3}$  cm/s by best fit of the calculated currents (Figure 4a,b, dashed lines) with the measured currents (Figure 4a,b, solid lines). It is noted that the sudden decrease in tip current at  $-0.6$  V in Figure 4a was also caused by hydrogen bubble formation and accumulation, similar to that shown in Figure 3a. The accuracy of our extracted kinetic information is also supported by a similar transient tip and substrate current measurements, as shown in Figure 4c,d. The comparisons of those tip and substrate currents for the V-MoS<sub>2</sub> and SV-MoS<sub>2</sub> substrates are shown in Figure S4a–d, which clearly show that the onset potential for SV-MoS<sub>2</sub> is about 0.1 V more positive than that of the V-MoS<sub>2</sub>, and the HER current of the SV-MoS<sub>2</sub> is about 4 times larger than that of V-MoS<sub>2</sub> at the same potential of  $-0.55$  V. Those comparisons are consistent with the previous observation that tensile strain further enhances the HER activity of the S vacancies.<sup>11</sup> The enhancement effect on HER by strain is more evidenced by comparing the extracted kinetic data, as summarized in Table 1. Though V-MoS<sub>2</sub> and SV-MoS<sub>2</sub> have similar formal potential  $E^0$  ( $-0.53$  V) and electron-transfer coefficient  $\alpha$  (0.4), the electron-transfer rate

**Table 1.** Extracted Kinetic Parameters for HER over V-MoS<sub>2</sub> and SV-MoS<sub>2</sub><sup>a</sup>

kinetic parameters	unstrained MoS <sub>2</sub> with S vacancies (V-MoS <sub>2</sub> )	strained MoS <sub>2</sub> with S vacancies (SV-MoS <sub>2</sub> )
$E^0$ (V vs Ag/AgCl)	$-0.53$	$-0.53$
$\alpha$	0.4	0.4
$k^0$ (cm/s)	$2.3 \times 10^{-4}$	$1.0 \times 10^{-3}$

<sup>a</sup>Other parameters used for fitting: proton diffusion coefficient of  $7.9 \times 10^{-5}$  cm<sup>2</sup>/s, hydrogen diffusion coefficient of  $5 \times 10^{-5}$  cm<sup>2</sup>/s, and proton concentration of 0.1 M.

constant  $k^0$  for SV-MoS<sub>2</sub> is about 4 times higher than that of V-MoS<sub>2</sub>.

Finally, the reliability of our SECM measurement was confirmed by four more repeated measurements of the V-MoS<sub>2</sub> and SV-MoS<sub>2</sub> regions on the same substrate (Figure S5, where the error bars represent the standard deviation). The very small errors suggest that our SECM measurements are reliable. We further confirmed this by complementary measurements using a three-electrode electrochemical compression cell that has a diameter of 3 mm (see our previous work for more details<sup>11</sup>). The average LSV ranges of both V-MoS<sub>2</sub> and SV-MoS<sub>2</sub> samples (eight samples total) are represented by shaded regions in Figure S6. One can see that the separate measurements agree reasonably well despite the different electrolytes and sample dimensions of the two distinct measurements.

## CONCLUSION

In summary, we have successfully applied the combinational approach of SECM measurements and multiphysics modeling to investigate the HER kinetics of newly identified active sites that consist of S vacancies on the basal plane of MoS<sub>2</sub>. The HER kinetic data for S vacancies were determined to be the formal potential  $E_v^0 = -0.53 \text{ V}_{\text{Ag}/\text{AgCl}}$ , electron-transfer coefficient  $\alpha_v = 0.4$ , and electron-transfer rate constant  $k_v^0 = 2.3 \times 10^{-4} \text{ cm/s}$ . In addition, the effect of tensile strain on the HER kinetics of S vacancies was quantified. About 2% uniaxial tensile strain increases the electron-transfer rate constant by almost 4-fold ( $k_{sv}^0 = 1.0 \times 10^{-3} \text{ cm/s}$ ), confirming that the tensile elastic strain indeed accelerates the HER kinetics of S vacancies in MoS<sub>2</sub>. The methodology presented here provides a general way to study the electrochemical kinetics of MoS<sub>2</sub>-like two-dimensional catalytic systems.

## ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.6b01377.

Supporting figures of the Raman spectra of strained and unstrained MoS<sub>2</sub>, optical image of a typical SECM sample, schematic SECM measurement setup, measured and calculated SECM approach curves, comparison of HER activities of V-MoS<sub>2</sub> and SV-MoS<sub>2</sub>, comparison of LSV curves of Pt, V-MoS<sub>2</sub>, and SV-MoS<sub>2</sub>, complementary LSV measurements using an electrochemical compression cell, simulated LSV curves at various  $k^0$ ; Supporting Note 1 of sample fabrication and characterization; Supporting Note 2 of substrate electrode preparation, Pt UME tip preparation, and SECM measurements; Supporting Note 3 of multiphysics simulation (PDF)

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### Notes

The authors declare no competing financial interest.

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